

# The Diffusion Current in Shallow Co-silicided Junctions

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**ABSTRACT** – This paper demonstrates that the diffusion current in shallow cobalt-silicided n<sup>+</sup>p junction diodes strongly increases with increasing deposited cobalt (Co) thickness ( $t_{Co}$ ). The increase is related to the enhanced minority carrier recombination current in the neutral emitter region which becomes smaller for higher  $t_{Co}$ . This is supported by the recombination lifetimes in the p-type substrate, derived from cross-sectional Microwave Absorption measurements, showing no dependence on the processing conditions. From the measured diffusion current density, one can derive the Gummel number of the emitter, which can be considered as a measure for the effective recombination properties in the highly doped region.

**KEY WORDS** - diffusion current, junction diode, Gummel number

บทคัดย่อ – ในบทความนี้จะแสดงให้เห็นว่ากระแสที่เกิดจากการแพร่ในไดโอดแบบรอยต่อ n<sup>+</sup>-p ที่มีโคบอลต์ซิลิไซด์ จะเพิ่มขึ้นอย่างมากเมื่อเพิ่มความหนาของโคบอลต์ที่เคลือบ ซึ่งการเพิ่มขึ้นสัมพันธ์กับการเพิ่มขึ้นของกระแสเนื่องจากการรวมตัวของพาหะส่วนน้อยในบริเวณเป็นกลางของอิมิตเตอร์ โดยบริเวณนี้จะมีขนาดเล็กลงเมื่อความหนาของโคบอลต์มีค่ามากขึ้น สมมติฐานนี้ได้รับการยืนยันจากค่าเวลาการรวมตัวของพาหะส่วนน้อยในฐานะรองชนิดที่ ซึ่งทำการวัดโดยการดูดกลืนคลื่นไมโครเวฟของภาคตัดขวาง (cross-sectional Microwave Absorption measurements) พบว่าค่าที่วัดได้ไม่ขึ้นอยู่กับเงื่อนไขของกระบวนการผลิต จากค่าความหนาแน่นของกระแสแพร่ที่วัดได้จะสามารถคำนวณค่าตัวเลขกัมเมลของชั้นอิมิตเตอร์ ซึ่งสามารถพิจารณาได้เหมือนกับ การวัดคุณสมบัติของการรวมตัวของพาหะส่วนน้อยในบริเวณที่มีการเจือสารสูงๆ

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## 1. Introduction

Downscaling CMOS technologies not only requires a reduction of the lateral device dimensions, but also of the source/drain junction depths. To achieve this goal, one has to reduce the thermal budget. However, this does not guarantee the complete removal of the ion implantation induced damage, which can enhance the diode leakage current when present in the depletion region. Another problem is related to the silicidation of the junction, which consumes part of the highly doped emitter region. It is well-known from the past that when the silicide interface comes too close to the junction, a drastic increase of the diode leakage current is observed [1-2]. It is therefore not a surprise that the formation of low-leakage shallow junctions is an interesting topic [3-5].

The focus in this paper is on the diffusion current density  $J_{diff}$  in shallow cobalt-silicided junctions. For deep p-n junctions, Shockley-Read-Hall theory predicts that the saturation value is given by [6]:

$$J_{diff} = \frac{qn_i^2}{N_A} \sqrt{\frac{D_n}{\tau_n}} + \frac{qn_i^2}{N_D} \sqrt{\frac{D_p}{\tau_p}} \quad (1)$$

with  $n_i$  the intrinsic carrier concentration;  $q$  the elementary charge;  $N_A$ ,  $N_D$  the doping density in the p- and n-side of the junction, respectively;  $D_n$  the diffusion coefficient of electrons (p-side) and  $D_p$  for holes (n-side) and  $\tau_p$  ( $\tau_n$ ) the corresponding minority carrier lifetime. For asymmetric junctions, i.e.  $N_A \ll N_D$ , the recombination current in the

neutral part of the highly doped  $n^+$  region (emitter) can be neglected compared with the diffusion current of the p-base (substrate). In such a case, the diffusion current density can be utilised to derive the minority carrier lifetime in the silicon substrate [7-8]. However, when the base (or substrate) is more highly doped, both terms in Eq. (1) need to be considered. In that case, it is possible to derive also the minority carrier lifetime in the emitter region, from the experimental diffusion current [6,9].

For shallow  $n^+p$  junctions, Eq. (1) can be approximated in first instance by [10]:

$$J_{diff} = \frac{qn_i^2}{N_A} \sqrt{\frac{D_n}{\tau_n}} + \frac{qn_i^2}{N_D} \frac{D_p}{d_{ne}} \quad (2)$$

which neglects surface recombination and assumes homogeneous material parameters in the emitter region. Here,  $d_{ne}$  is the thickness of the neutral emitter region, extending from the silicide interface to the edge of the depletion region on the  $n^+$  side. As can be seen from Eq. (2), for sufficiently small  $d_{ne}$ , the recombination current in the highly doped  $n^+$  region may no longer be negligible compared with the substrate contribution. This can, for example, occur when during silicidation a large part of the highly doped region is consumed. This effect is studied here for a series of Co silicided  $n^+p$  junctions, corresponding to different thickness of deposited Co ( $t_{Co}$ ).

## 2. Experiments

Co-silicided  $n^+p$  junctions were processed on 150 mm p-type Cz wafers, with a typical doping density of  $6 \times 10^{14} \text{ cm}^{-3}$ , as derived from high-frequency Capacitance-Voltage (C-V) measurements. The junctions have been formed by an arsenic implantation at 70 keV, using a dose of  $3 \times 10^{15} \text{ cm}^{-2}$ . Dopant activation was achieved by an 1100 °C rapid thermal anneal (RTA) for 10 sec. The junctions have been silicided, using a different thickness of the sputtered Co layer  $t_{Co}$ , as indicated in Table I. Figure 1 shows cross-sectional Transmission Electron Microscopy (TEM) micrograph for the Co thickness of 30 nm. The non-silicided reference wafer received a standard Al metallization, corresponding to an original junction depth ( $d_j$ ) of about 0.15  $\mu\text{m}$ . The arsenic profile in the  $n^+$  region has been determined by SIMS measurements. The remaining emitter thickness of Table I is calculated assuming a resulting silicide thickness of:

$$t_e = d_j - 3.73t_{Co} \quad (\text{in nm}) \quad (3)$$

Current-Voltage (I-V) and C-V measurements have been performed at 25 °C on five large area ( $A=10 \text{ mm}^2$ ) and five large perimeter diodes ( $P=8.04 \text{ cm}$ ) per wafer. Combining either the forward or the reverse I-V characteristics of the area and perimeter diodes, one can separate first the geometrical leakage current components [11]. Next, the saturation value of the diffusion current density can be extracted either from the forward I-V, as demonstrated previously [8], or from the reverse volume current density, using a recently developed procedure, which is also valid for shallow junctions, suffering from field-enhanced generation effects [12]. Combining the C-V data of the area and perimeter diode yields the doping profile  $N_A(W)$  in the substrate and the depletion width ( $W$ ), which are used in the current separation [8,11-12].

Table I. Processing parameters and ideality factor ( $m$ ) of the exponential part of the forward diode characteristics.

deposited Co thickness ( $t_{Co}$ ) (nm)	emitter thickness ( $t_e$ ) (nm)	ideality factor ( $m$ )
0 (no)	150	1.10
12	106	1.14
20	77	1.12
30	41	1.03
20 + 10 nm Ti	77	1.06

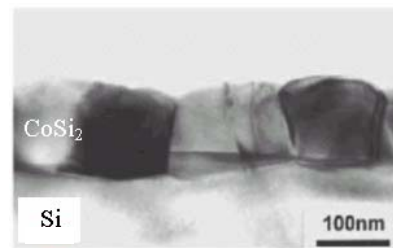


Figure 1. Cross-sectional Transmission Electron Microscopy (TEM) micrograph for the Co thickness of 30 nm.

The recombination lifetime in the bulk has also been measured directly using time-resolved microwave (MW) photoconductivity measurements at Vilnius University. The principle of the technique and the data analysis is described elsewhere [13].

### 3. Results

The silicidation has a very strong impact on the I-V characteristics as can be derived from Fig. 2. The current in both forward and reverse polarization increases strongly with increasing Co thickness, in agreement with previous studies [1-3]. Nevertheless, the ideality factor (Table I) of the diodes are good, ranging between 1.03 and 1.14, in agreement with other reports [3]. This indicates that in the ideal forward region, the current is dominated by the diffusion mechanism. No important generation-recombination (GR) or tunneling component is observed there.

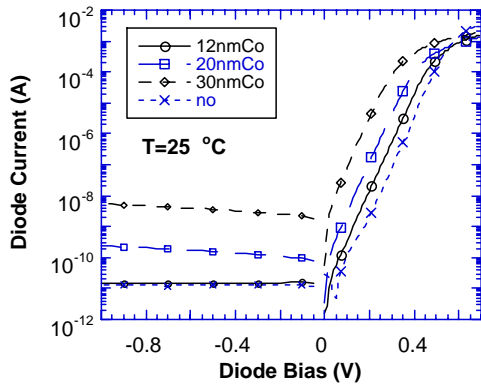


Figure 2. I-V characteristic at 25 °C for a set of 10 mm<sup>2</sup> diodes, corresponding to different deposited Co thickness.

Figure 3 represents the saturation value of the diffusion current density extracted from the forward and the reverse I-V. Note first of all that within 20 % the same value for  $J_{diff}$  is derived from both methods, demonstrating the consistency of the extraction [8,12]. Secondly,  $J_{diff}$  shows a strong, close to exponential variation with  $t_{Co}$  (or with  $t_e$ , taking account of Eq. (3)). The observed increase is much stronger than the reciprocal one predicted by Eq. (2), suggesting that this equation is rather qualitative than quantitative.

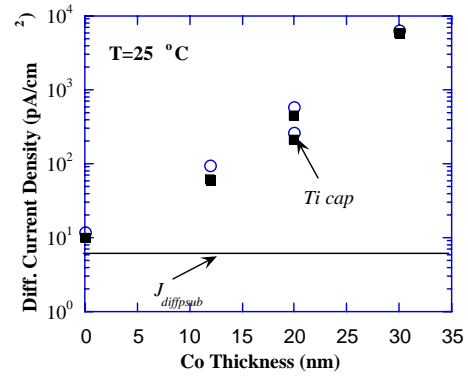


Figure 3. Diffusion current density derived from the forward (open circles) and the reverse (closed squares) I-V, for the diode set of Table I.

It has been suggested on a number of occasions that silicidation injects point defects in the silicon and therefore introduces generation/recombination centres in the diode depletion region and deeper in the substrate, which increase the generation and the diffusion current [1-2]. In order to test this hypothesis, it was decided to perform MW measurements, the results of which are summarized in Fig. 4. From the effective recombination lifetime derived from the transient microwave signal, following a laser pulse excitation, both the bulk lifetime and the surface recombination velocity can be extracted [13]. Here, the light excitation was performed in a cross section of the sample, enabling the assessment of the lifetime parameters close to the junction, which was not removed for the analysis. It is evident from Fig. 4 that the silicidation has a marginal impact on the bulk recombination lifetime, which is around 30  $\mu$ s.

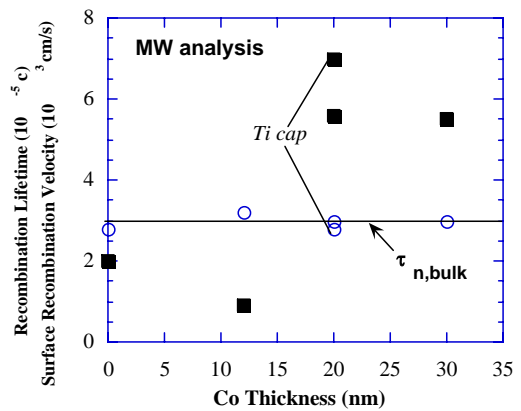


Figure 4. MW bulk recombination lifetime (circles) and surface recombination velocities (squares) as a function of the Co thickness.

Having ruled out a possible change in the bulk recombination properties, we are led to the conclusion that the change in  $J_{diff}$  is related to the emitter region. However, as shown above, Eq. (2) is not an adequate description to account for the recombination in the neutral emitter. A better approximation is given by [3]:

$$J_{diff,n} = \frac{qn_i^2}{G} \quad (4)$$

with  $G$  the Gummel number defined by:

$$G = \int_0^{d_{ne}} \frac{N_D(x)}{D_{heff}(x)} \exp\left(\frac{\Delta E_G}{kT}\right) dx \quad (5)$$

whereby  $N_D(x)$  is the free carrier profile in the  $n^+$  region,  $D_{heff}$  the effective hole diffusivity and  $\Delta E_G$  the corresponding band-gap narrowing which is also position dependent. Equation (4), on the one hand, takes account of the high-doping effects on the band-gap ( $n_i$ ) and on the diffusivity, and, on the other hand, accounts for the profile in the doping density. Literature models have been used for  $\Delta E_G$  and  $D_{heff}$ , in order to calculate  $G$  [14].

Experimental  $G$ -values have been derived from the saturation diffusion current density as follows. First, for all diodes the diffusion current density corresponding to a lifetime of 30  $\mu$ s and a  $D_n = 35$   $\text{cm}^2/\text{s}$  has been calculated and subtracted from the measured value. The result should correspond to the emitter contribution, given by Eq. (4). Figure 5 compares the experimental with the calculated Gummel number. From the figure, a good qualitative agreement is observed, whereby both curves seem to be shifted over about 20 nm. More detailed calculations reveal that the Gummel number is very sensitive to the actual dopant profile in the emitter region close to the junction. However, it is clear that for shallow junctions below 100 nm, an increase in the diffusion current density by roughly one decade occurs for an emitter consumption of 30 nm. This is in line with other estimates which can be found in literature [5].

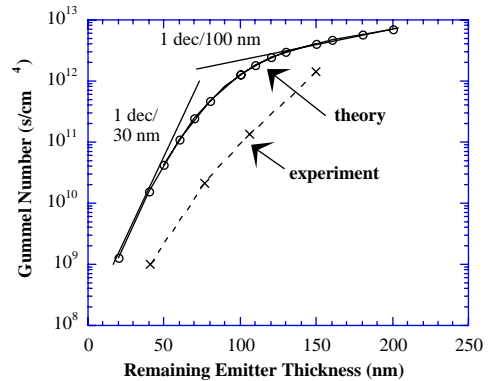


Figure 5. Calculated and experimental Gummel number for different starting junction depths versus the remaining emitter thickness.

From the calculations follows that the steep increase of the diffusion current density is explained by the fact that for shallow junctions, the arsenic profile is governed by the exponential tail of the profile, while for deeper junctions with respect to the silicide interface, the flat part of the dopant profile is not consumed and corresponds to an increase of 1 decade in diffusion current for every 100 nm consumed. This implies that the diffusion current may impose a limitation to the further down-scaling of p-n junctions, when it exceeds a certain threshold. It is furthermore a sensitive function of the activated dopant profile [5,10]. Of course, for very shallow junctions below about 50 nm, other transport mechanisms like thermionic emission become dominant and should be considered [10].

## 4. Conclusion

The impact of the diffusion current due to minority carrier transport in the remaining  $n^+$  emitter has been clarified. It is believed that it imposes a fundamental lower limit for the diode leakage, which should be considered for the development and modeling of (ultra) shallow junctions.

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